

Acta Cryst. (1978). **A34**, 830

A note on On the diffraction of X-rays by face-centred cubic crystals containing extrinsic stacking faults
by **C. J. Howard**. By YOSHITO TAKAKI, *Department of Physics, Osaka Kyoiku University, Tennoji, Osaka 543, Japan*
and JIRO KAKINOKI, † *Faculty of Engineering, Setsunan University, Ikeda-Naka, Neyagawa 572, Japan*

(Received 18 February 1978; accepted 12 May 1978)

The intensity of X-ray scattering from f.c.c. crystals containing extrinsic stacking faults with two fault parameters [see Howard (1977). *Acta Cryst.* **A33**, 29–32] has been calculated.

Howard (1977) has recently calculated the intensity of X-rays diffracted from an f.c.c. crystal containing extrinsic stacking faults by a method using difference equations (Holloway & Klamkin, 1969). He has also given a more general model for faulting with two fault parameters, p and q , but no solution has been given. In this note, the intensity formula for the above two-parameter model is derived by using the matrix method given by Kakinoki (1967).

† Permanent address: Higashi 1-17, Hagiwaradai, Kawanishi 666, Japan.

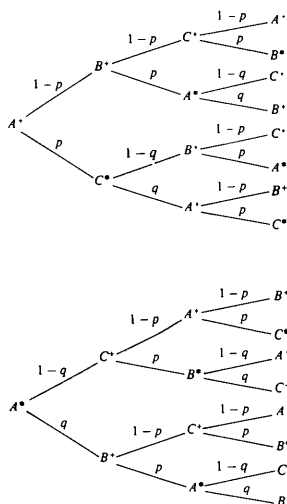


Fig. 1. The stacking sequence of layers for the two-parameter model given by Howard (1977). Analogous probability trees can be drawn from B^+ , B^* , C^+ and C^* .

Table 1. The 'P-table' representing the same stacking sequence as that shown in Fig. 1

	A^+	A^*	B^+	B^*	C^+	C^*
$w/3$	A^+		$1-p$			p
$w'/3$	A^*		q		$1-q$	
$w/3$	B^+	p			$1-p$	
$w'/3$	B^*	$1-q$			q	
$w/3$	C^+	$1-p$		p		
$w'/3$	C^*	q	$1-q$			

Fig. 1 shows the probability trees obtained on the basis of Howard's two-parameter model. According to Kakinoki (1967), the stacking sequence of layers shown in Fig. 1 can be represented by the 'P-table' as shown in Table 1, where $w = 1/(1+p)$ and $w' = p/(1+p)$, $w/3$ and $w'/3$ being probabilities of occurrence of layers in a crystal of types A^+ (B^+ or C^+) and A^* (B^* or C^*) respectively. The intensity formula for the diffuse scattering can then easily be derived as follows:

$$I = \frac{P_0 [P_1 - 2P_2 \cos \pi(l \mp \frac{1}{3})]}{Q_0 + Q_1 \cos \pi l \pm Q_2 \sin \pi l + Q_3 \cos 2\pi l \pm Q_4 \sin 2\pi l} \quad (1)$$

where the hexagonal coordinates $[hk.l]$ are the same as those of Howard; the upper and lower signs correspond to the cases $(h-k) \equiv 1 \pmod{3}$ and $(h-k) \equiv 2 \pmod{3}$ respectively; $P_0 = 3p(1-p+pq)/(1+p)$, $P_1 = 2-q-pq$, $P_2 = 1-q$, $Q_0 = 2(1-p+p^2) - 3p^2q(1-q)$, $Q_1 = (1-p)(1+2p-3pq)$, $Q_2 = -\sqrt{3}(1-p)(1+pq)$, $Q_3 = p(1-3q)$ and $Q_4 = -\sqrt{3}p(1-q)$.

For $q = p$, equation (1) reduces to

$$I = \frac{3p(1-p)(1-p+p^2)[2+p-2\cos\pi(l \mp \frac{1}{3})]/(1+p)}{A} \quad (2)$$

with

$$A = 2(1-p+p^2) - 3p^3(1-p) + (1-p)^2(1+3p) \\ \times \cos \pi l \mp \sqrt{3}(1-p)(1+p^2) \sin \pi l \\ + p(1-3p) \cos 2\pi l \mp \sqrt{3}p(1-p) \sin 2\pi l$$

which is identical with equation (14) of Howard, though these two equations have apparently quite different forms. For $q = 0$, an equation the same as that given by Johnson [1963; see equation (15) of Howard] can be obtained.

References

- HOLLOWAY, H. & KLAMKIN, M. S. (1969). *J. Appl. Phys.* **40**, 1681–1689.
HOWARD, C. J. (1977). *Acta Cryst.* **A33**, 29–32.
JOHNSON, C. A. (1963). *Acta Cryst.* **16**, 490–497.
KAKINOKI, J. (1967). *Acta Cryst.* **23**, 875–885.